

1 Imaging and modeling the ionospheric airglow
2 response over Hawaii to the tsunami generated by
3 the Tohoku Earthquake of 11 March 2011

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4 Although only centimeters in amplitude over the open ocean, tsunamis can
5 generate appreciable wave amplitudes in the upper atmosphere, including
6 the naturally occurring chemiluminescent airglow layers, due to the expo-
7 nential decrease in density with altitude. Here, we present the first obser-
8 vation of the airglow tsunami signature, resulting from the 11 March 2011
9 Tohoku earthquake off the eastern coast of Japan. These images are taken
10 using a wide-angle camera system located at the top of the Haleakala Vol-

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11 cano on Maui, Hawaii. They are correlated with GPS measurements of the
12 total electron content from Hawaii GPS stations and the Jason-1 satellite.
13 We find waves propagating in the airglow layer from the direction of the earth-
14 quake epicenter with a velocity that matches that of the ocean tsunami. The
15 first ionospheric signature precedes the modeled ocean tsunami generated
16 by the main shock by approximately one hour. These results demonstrate
17 the utility of monitoring the Earths airglow layers for tsunami detection and
18 early warning.

1. Introduction

19 As a tsunami propagates over the open ocean, atmospheric gravity waves (AGWs)
20 forced by the centimeter-level surface undulations are generated [*Hines, 1960; Peltier and*
21 *Hines, 1976*]. Due to the decrease in atmospheric density and the requirement that wave
22 momentum is conserved, the amplitudes of the forced AGWs can reach several kilometers
23 at altitudes where the neutral atmosphere coexists with the plasma in the ionosphere
24 (~ 250 km), perturbing the ionosphere through collisions between the neutrals and the
25 charged particles [*Occhipinti et al., 2008; Hickey et al., 2009*]. This ionospheric response
26 was originally conjectured in the 1970s and formed the basis of a proposed tsunami early
27 detection technique [*Najita et al., 1974*]. In the intervening decades, little progress was
28 made in utilizing any ionospheric observation methods due to the relative paucity of
29 required measurements.

30 The first observation using the very dense Japanese GPS Earth Observation Network
31 (GEONET) was reported in 2005, with the imaging of the ionospheric perturbations re-
32 lated to the trans-Pacific tsunami generated by the Mw=8.2 earthquake that occurred
33 in Peru on 23 June 2001 [*Artru et al., 2005*]. Additionally, several observations were
34 performed during the 2004 Sumatra tsunami. These observations rely on estimates of
35 the total electron content (TEC), the integrated electron density between a specific GPS
36 satellite and receiver [*Liu et al., 2006; Lognonné et al., 2006*] or between a satellite-based
37 altimeter and the sea surface, and have been numerically reproduced [*Occhipinti et al.,*
38 *2006*]. Subsequent studies have shown systematic ionospheric tsunami detection utiliz-
39 ing GPS-derived TEC was possible using smaller networks, such as on Hawaii [*Rolland*

40 *et al.*, 2010]. However, as the density and coverage of the available GPS network decrease,
41 the ability to “image” the ionospheric response is lost. Furthermore, the ability to de-
42 tect the tsunami is dependent on the constantly changing geometry of the GPS satellite
43 constellation in addition to the background ionospheric structure.

44 Here and for the first time, we utilize a highly sensitive, wide-angle camera system to
45 image the tsunami-driven ionospheric response to the 11 March 2011 Tohoku earthquake.
46 From a single instrument located on the Haleakala Volcano on Maui, Hawaii, we are able
47 to image a 10^6 km² region of the ionosphere at high spatial ($\sim 1 - 5$ km, elevation angle
48 dependent) and temporal (~ 5 min) resolutions. This is done by observing the airglow
49 layer at approximately 250 km in altitude caused by the dissociative recombination of
50 O_2^+ [Link and Cogger, 1988], which emits photons at 630.0 nm. Modeling studies have
51 suggested that appreciable modulations in the 630.0-nm intensity should be caused by
52 tsunami-driven AGWs [Hickey *et al.*, 2010]. However, until now, this has never been
53 demonstrated.

2. Data Presentation and Analysis

54 As reported by the United States Geological Survey (USGS), the Mw=9.0 Tohoku
55 earthquake occurred at 05:46:23 UT off the east coast of Honshu, Japan (38.322° N,
56 142.369° E). Based on the measurement of the Deep-ocean Assessment and Reporting
57 of Tsunamis (DART) buoy 51407 located near the Big Island of Hawaii (19.6° N, 203.5°
58 E), the arrival of the tsunami was predicted at 13:07 (PTWC message 5) and the first
59 maximum amplitude (15 cm) is reported at 13:37 UT (PTWC message 27). Tsunami
60 Travel Time software [Wessel, 2009] using the USGS seismic source predicts an arrival

61 time of 13:17 UT at the DART 51407. The imaging system utilized in this study began
62 routine observations at 10:32 UT and continued observing until 15:19 UT. The skies were
63 predominately clear, with only a few low-lying clouds observed over the course of the night
64 and we can therefore be confident structures observed are actually in the airglow layer.

65 Initial analysis of the raw images collected of the 630.0-nm emission on this night
66 exhibited little of interest related to the tsunami. Typical structures observed in this
67 imaging system can include north-south aligned and eastward propagating equatorial
68 plasma bubbles [e.g., *Kelley et al.*, 2002; *Makela et al.*, 2004], northwest-southeast aligned
69 and southwestward propagating medium-scale traveling ionospheric disturbances [MSTID;
70 e.g., *Makela et al.*, 2009; *Miller et al.*, 2009], and poleward propagating brightness waves.
71 A very faint MSTID is observed from the beginning of the observations until approxi-
72 mately 11:30 UT (much earlier than the expected tsunami arrival time) and a distinct
73 brightness wave is observed from 12:45-13:45 UT. Neither of these features are tsunami
74 related.

75 A more detailed analysis was performed in which the images were projected onto an
76 assumed emission layer altitude of 250 km and then passed through a filter to isolate
77 specific wave periods. Three different length-8 FIR filters were used with passbands of 0.3-
78 1.0 mHz, 1.0-1.7 mHz, and 0.3-1.7 mHz. The imaging data were sampled at approximately
79 3.6 mHz. A Gabor filter was then used to estimate the spatial frequency and propagation
80 angle of waves within each passband. This filtering revealed both a long-period wave
81 ($T = 26.2 \pm 3.1$ min; $\lambda = 290.0 \pm 12.5$ km; $v = 184.5 \pm 33.8$ m/s) and a short-period
82 wave ($T = 14.2 \pm 2.7$ min; $\lambda = 189.9 \pm 4.9$ km; $v = 222.9 \pm 52.4$ m/s) propagating in the

83 direction of the tsunami as it passed by the Hawaiian Islands. Example filtered images
84 are shown in Figure 1 and a movie showing all of the filtered images collected on this
85 night is included in auxiliary online material.

86 The long- and short-period waves have slightly different arrival directions ($132^\circ \pm 1^\circ$
87 and $136^\circ \pm 1^\circ$ azimuth, respectively) before approximately 13 UT, suggesting that the
88 source of these two wave packets may be different. However, ray tracing of the tsunami
89 (not shown) indicates that the two azimuths reported for the ionospheric early waves
90 originate from the rupture zone reported by USGS. This focusing can be attributed to the
91 Hawaiian-Emperor seamount chain bathymetry, which acts as a tsunami waveguide. After
92 approximately 13 UT, both waves show an arrival direction of $134^\circ \pm 1^\circ$. The alignment
93 and propagation direction of the observed waves is inconsistent with the other types of
94 structures typically seen in this imaging system and their occurrence at the same time,
95 direction, and speed as the tsunami makes it reasonable to attribute these airglow waves
96 to the tsunami propagation.

97 To confirm this hypothesis, we have obtained GPS 30-sec sampled data from 52 re-
98 ceivers on the Hawaiian Islands available on the UNAVCO public ftp website ([ftp://data-
99 out.unavco.org/](ftp://data-out.unavco.org/)). Following the methodology of [Rolland *et al.*, 2010], we have extracted
100 the tsunami signature in these data. Similarly, we have obtained TEC data measured
101 with the dual-frequency altimeter on the Jason-1 satellite and have extracted the tsunami
102 signature as done previously for the Sumatra tsunami [Occhipinti *et al.*, 2006]. We have
103 compared these measurements to the imaging data collected on this night and the waves

104 seen in the TEC data are collocated in space and time with the waves observed in the
105 images, as shown in the movie.

106 Time of arrival diagrams constructed using the airglow data and GPS-derived TEC
107 data are presented in Figure 2. In this representation, structures that are moving with
108 the tsunami (in terms of both direction and speed) appear as vertical bands. Structures
109 arriving before the tsunami appear at negative time values on the x axis, while structures
110 following the tsunami appear at positive time values. It is evident from this figure that
111 structures are seen in the ionosphere approximately one hour ahead of the tsunami.

3. Discussion

112 The observed early waves, if indeed related to the earthquake, may have several com-
113 peting origins. They might be caused by (a) an infrasonic wave generated at the earth-
114 quake source propagating in the ionosphere faster than the ocean tsunami wave front, (b)
115 a plasma diffusion oscillation related to the sea-level disturbances of the tsunami wave
116 front, or (c) the initial tsunami wave generated by a pre-rupture processes.

117 If the observed early waves were caused by an infrasonic wave generated at the earth-
118 quake source and traveling in the ionosphere, we would expect them to be observed prop-
119 agating from the geometrical direction between the source region and Hawaii. This would
120 be an azimuth angle of approximately 123° which differs significantly from the observed
121 azimuth angles. Secondly, if this were the case, we would not expect to see any associated
122 early perturbations at the ocean surface which, as shown below, are observed. Neither
123 such an infrasonic wave nor the ocean floor deformation generated by the tsunami loading
124 can generate the observed waves.

125 To examine the second potential cause of these early waves, we use a sea-level modeling
126 of the tsunami as generated following the methodology of *Hébert et al.* [2007] and using
127 the USGS finite fault model (provided at <http://earthquake.usgs.gov/earthquakes>) for
128 the Tohoku earthquake. We model the resultant AGWs and associated ionospheric waves
129 with two different methods, assuming either a pure gravity wave regime [*Occhipinti et al.*,
130 2006, 2008] or a gravity-acoustic regime [*Kherani et al.*, 2009]. A comparison of the
131 observed ionospheric waves in the airglow and each model's response at 250 km (the
132 altitude of the assumed peak in the airglow intensity) is presented in Figure 3. The
133 agreement between the observations of waves after the arrival of the tsunami and modeling
134 in both cases is not only conclusive in arrival time and wave structure azimuth, but also
135 for several features appearing in the wave structures [*Occhipinti et al.*, submitted, 2011],
136 which are observed about one hour after the passage of the tsunami. However, no evidence
137 of the early wave is seen in either model.

138 To test the third hypothesis, a closer analysis of the DART data has been performed.
139 The large tidal signature in the data was removed using the 45-day average tidal signature
140 from the DART measurements on 11 March 2011 using a least squares fitting process. The
141 raw data and the data with the tide removed in this way are presented in the top two
142 panels of Figure 4. The residual sea level measurement shows a 1-cm retreat preceded
143 by a slow ramp extending approximately 90 minutes before the arrival of the tsunami
144 at the DART buoy. Superposed on this ramp are higher-frequency waves. Additional,
145 independent methods for removing the tidal signature, including a high-pass filter and a
146 Fourier methodology, were performed and confirm the presence of the 1-cm retreat and

147 preceding ramp with superposed waves features, as shown in the third panel of Figure 4.
148 A spectral analysis of the residual signal, presented in the bottom-right panel of Figure 4,
149 indicates that the energy in these early waves has a dominant period of about 28 min with
150 lesser contributions at shorter periods, which include the 14 min period. This suggests
151 that the early waves observed in the ionosphere are related to these small oscillations seen
152 on the ocean surface.

153 A model of the expected sea level variations, taking into account the characteristics of
154 the fault source reported by the USGS, is presented as the red line in the second panel
155 of Figure 4. A second, independent model of the expected sea level variations was also
156 run following the method of [*Hébert et al.*, 2007], and is presented as the red line in the
157 third panel of Figure 4. Neither modeling methodology generates the observed ramp
158 and associated small fluctuations seen in the residual signal. Analysis of other DARTs
159 (43412, 51406, 43413 and, with more noise, 52406 and 32413) indicate similar features.
160 As our ocean-to-atmosphere modeling presented above, which relies on the sea-surface
161 models, does not generate the observed early ionospheric waves and since the infrasound
162 hypothesis would not generate such a sea level displacement, we conclude that a seismic
163 origin for this early wave, possibly due to pre-rupture processes and not included in the
164 USGS fault source, is likely.

165 As shown here, the use of an imaging system observing the airglow layers in the iono-
166 sphere provides a powerful tool for monitoring the passage of tsunamis. Dense networks
167 of instruments are not required to produce an image, as is the case using GPS-derived
168 TEC, and classical image processing techniques can be applied to enhance the utility of

169 the information derived from the measurements. However, the airglow technique does
170 have its drawbacks. Specifically, the ground-based camera system utilized here is con-
171 strained to taking images when both the sun and the moon are below the local horizon
172 and the skies overhead are clear, conditions that will not be met for every tsunami event.
173 Furthermore, these types of systems are constrained to land-based locations making mon-
174 itoring the entire Pacific Ocean, for example, impossible. However, an imaging system on
175 a space-based platform would not suffer from these limitations. Several emissions sensitive
176 to the ionospheric density are observable from space on both the night- and dayside, with
177 the most promising being the far ultraviolet OI emission at 135.6 nm. The response of
178 this emission to a tsunami-induced AGW has been modeled in *Hickey et al.* [2010]. A
179 geostationary platform would allow for continual monitoring of large regions of the Earth
180 and usher in a new era of tsunami monitoring.

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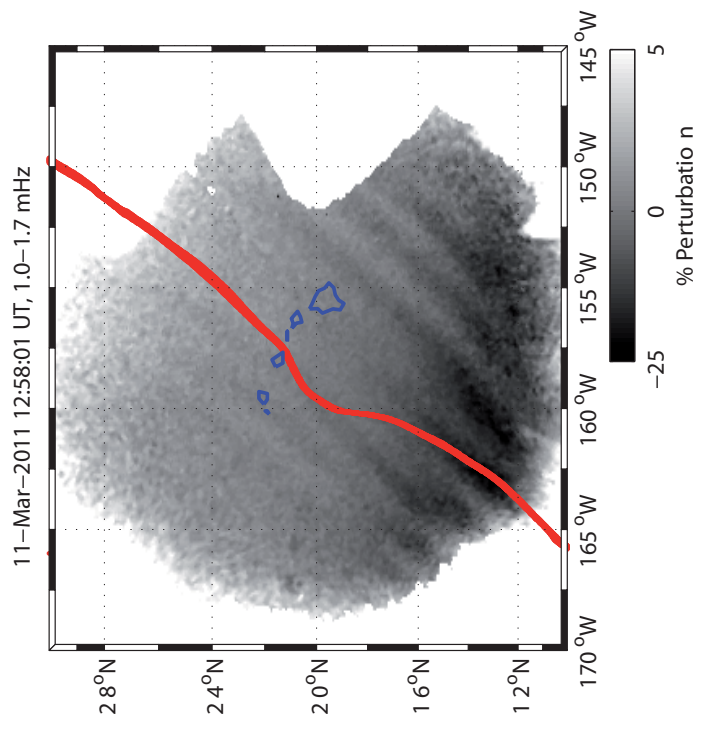
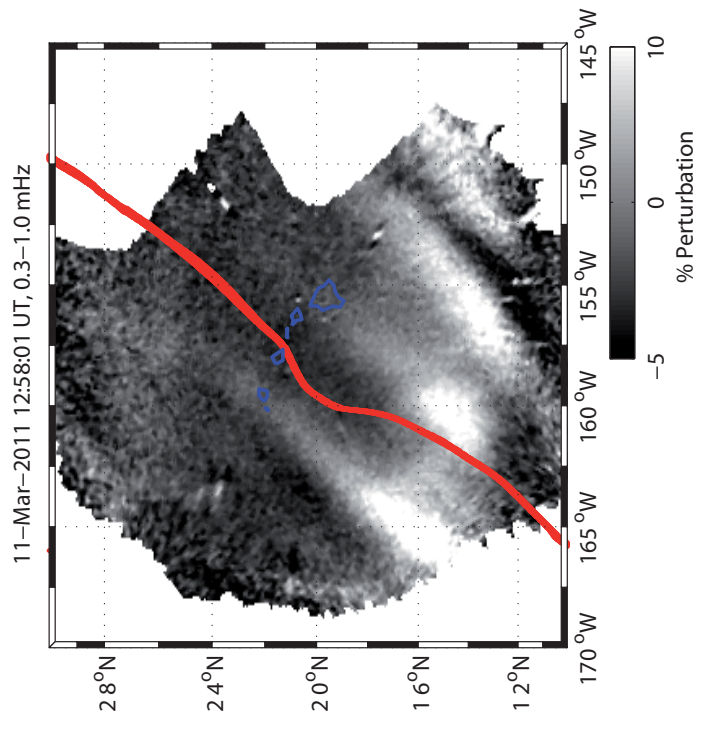
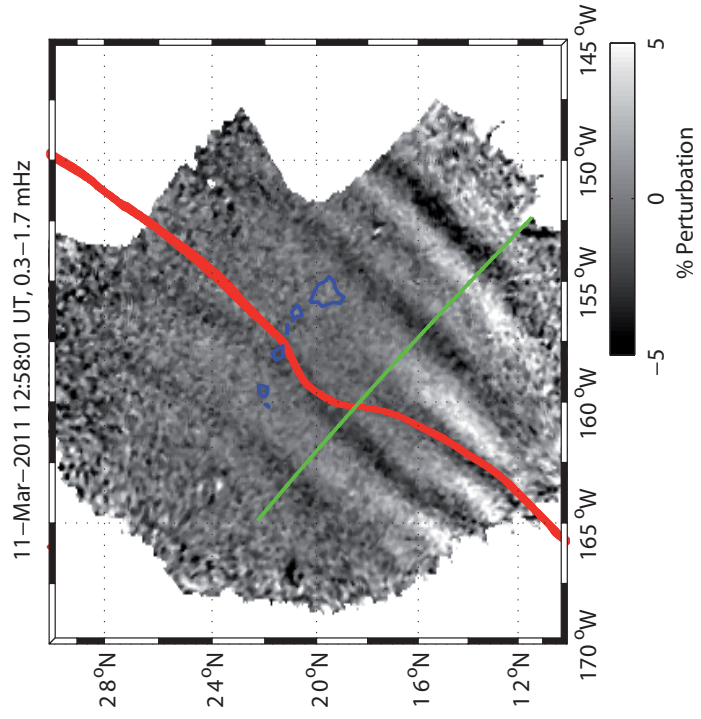
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Figure 1. Example of 630.0-nm images processed using length-8 FIR filters with passbands of (left) 0.3-1.7 mHz, (middle) 0.3-1.0 mHz to highlight the 26.2-min period waves, and (right) 1.0-1.7 mHz to highlight the 14.2-min period waves. The red line in each image indicates the tsunami location at the time of the image. The green line in the left figure indicates the line from which intensities were taken to construct Figure 2.

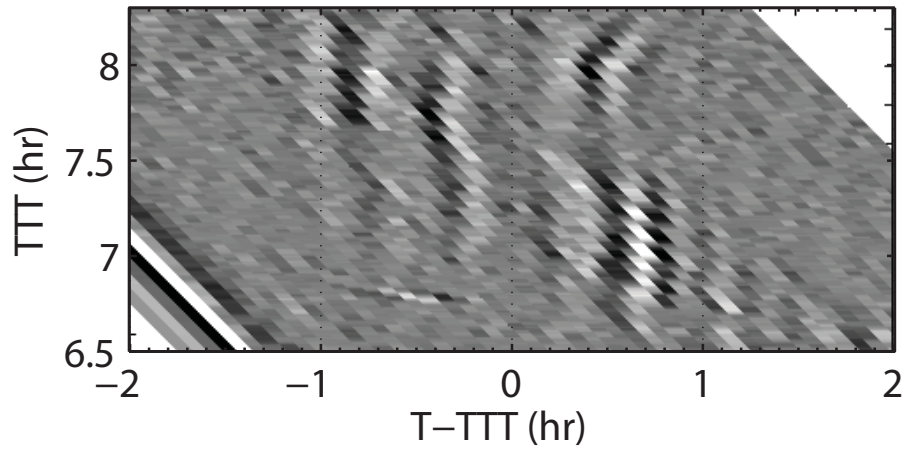
Figure 2. Travel-time diagrams for the (top) 630.0-nm emission intensity and (bottom) vertical TEC. The intensities for the top figure were extracted along the green line shown in Figure 1. Structures that are moving at the same speed and direction as the tsunami appear as vertical bands. Negative values along the x-axis indicate structures that appear before the arrival of the tsunami. Both panels are for the respective datasets temporally filtered between 0.3-1.7 mHz.

Figure 3. Comparison of (left) differenced 630.0-nm emission intensity observed at 13:20 and 13:22 UT from Hawaii, (center) electron density at 250 km from a gravity-acoustic model [Kherani *et al.*, 2009], and (right) normalized vertical wind velocity at 250 km from a pure gravity wave model [Occhipinti *et al.*, 2006, 2008, submitted, 2011]. In each case, the red line indicates the tsunami location at the time of the image.

Figure 4. (First panel) Tidal-detrended data (blue line) from DART 51407 and PREM synthetics (red line) centered at the theoretical tsunami arrival time (13:17 UT). (Second panel) Same as the first panel, but zoomed in around the tsunami arrival time. In both panels, the tidal-detrended results are obtained using a least-square fitting over a 45-day time series. The PREM synthetics are computed from source inversion of surface waves in the band 250-450 sec. The synthetics in the first panel are shifted by +0.025 m to allow for a clearer presentation of the data and comparison with the surface waves observed at the DART station about 420 min before the tsunami arrival. (Third panel) Residual sea level variations recorded on DART 51407 filtered from the oceanic tide using high-pass filtering (blue line) or a trend computed from a Fourier model (green line) compared to those using a model based on the co-seismic initial deformation without any precursor signal (red line). The power spectrogram of the data (bottom figure right) indicating a dominant frequency of 0.6 mHz (27.77 min) starting to emerge from the background noise about 20 min prior to the tsunami arrival time. The power spectrogram at the bottom left corresponds to the DART data, with a 5-min cosine taper starting after the tsunami arrival time and with PREM synthetics added before demonstrating that the surface waves are not responsible for generating the long-period waves.



630.0-nm TTT; 0.3–1.7 mHz; 11-Mar-2011



TEC TTT; 0.3–1.7 mHz; 11-Mar-2011

